

## Dynamics of Vortex Line in Presence of Stationary Vortex

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We study dynamics of infinitely thin vortex line of zero vorticity and unit local induction parameter in a velocity field created by infinitely thin stationary vortex line of vorticity  $q$ . If the shape of line is a one-valued function of vertical coordinate  $z$ , position of the moving vortex line is described by a complex function

$$\Psi = \Psi(z, t) \quad \Psi = x + iy$$

Function  $\Psi(z, t)$  satisfies equation

$$\frac{\partial \Psi}{\partial t} = i \left( \frac{\partial}{\partial z} \frac{1}{\sqrt{1 + |\Psi_z|^2}} \Psi_z + \frac{q}{\Psi} \right) \quad (1)$$

If  $q > 0$ , both vortices rotate in the same direction (co-rotating case); if  $q < 0$ , vortices rotate in the opposite direction (anti-rotating case). If  $q = 0$ , the moving vortex is free and Equation (1) is an integrable system equivalent to the focusing Nonlinear Schrodinger equation (NLSE).

Equation (1) is Hamiltonian

$$\frac{\partial \Psi}{\partial t} = i \frac{\delta H}{\delta \bar{\Psi}} \quad H = \int \left\{ 2(R - 1) - q \ln \frac{|\Psi|^2}{|\Psi_0|^2} \right\} dz \quad (2)$$

Here  $R = \sqrt{1 + |\Psi_x|^2}$ . We assume that  $|\Psi| \rightarrow |\Psi_0|^2$  at  $z \rightarrow \infty$ . Besides  $H$ , it has the following constants of motion

$$N = \int (|\Psi|^2 - |\Psi_0|^2) dz$$

$$P = i \int |\Psi \bar{\Psi}' - \bar{\Psi} \Psi'| dz$$

Equation (1) has helix-type exact solutions

$$\begin{aligned} \Psi &= A e^{ikz - i\omega t} \\ \Omega &= \frac{k^2}{\sqrt{1 + A^2 k^2}} - \frac{q}{A^2} \end{aligned} \quad (3)$$

Study of spiral stability with respect to small perturbation

$$\Psi = Ae^{ikz-i\omega t} \left( 1 + \delta Ae^{-i\omega t+ipz} \right)$$

leads to the following dispersion relation

$$\omega^2 = \frac{p^2}{1+k^2A^2} \left\{ \frac{p^2}{1+k^2A^2} - \frac{k^4A^2}{1+k^2A^2} + \frac{2q}{A^2} \right\} \quad (4)$$

If  $q = 0$ , the spiral is unstable. This is nothing but modulational instability in the focusing NLSE. Also, the helix is unstable in the anti-rotating case  $q < 0$ . In the co-rotating case, the spiral is stable if

$$\frac{2q}{A^2} > \frac{k^4A^4}{1+k^2A^2}$$

In particular, co-rotating straight vortex  $k = 0$  is stable. If  $q < 0$ , such vortex is unstable; this is an analog of well-known Crow instability of a pair of anti-parallel vortices.

In the anti-rotating case, the development of helix instability leads to gluing of both vortices in a finite time. This process can be treated as reconnection of vortices. The final stage of this reconnection can be described by self-similar solution of equation (1):

$$\Psi = (t_0 - t)^{1/2+i\nu(q)} \Phi \left( \frac{z}{\sqrt{t_0 - t}} \right) \quad (5)$$

Here  $\nu(q)$  is a nonlinear eigenvalue to be found numerically.

In the co-rotating case, Equation (1) has a rich family of solitonic solutions, which stabilize at  $z \rightarrow \infty$  to a stable spiral. Such solitonic solution can be sophisticated, in particular, knotted. Their stability is not properly studied yet.